

## ULTRAVIOLET, X-RAY, AND INFRARED OBSERVATIONS OF HDE 226868 = CYGNUS X-1

A. TREVES AND L. CHIAPPETTI  
 Istituto di Fisica dell'Università, Milan, Italy

E. G. TANZI  
 Laboratoria de Fisica Cosmica e Tecnologie Relative, Milan, Italy

M. TARENGHI  
 European Southern Observatory

H. GURSKY, A. K. DUPREE, L. W. HARTMANN, J. RAYMOND, R. J. DAVIS, AND J. BLACK  
 Harvard-Smithsonian Center for Astrophysics

T. A. MATILSKY  
 Department of Physics, Rutgers University

P. VANDEN BOUT AND F. SANNER  
 University of Texas at Austin

G. POLLARD AND P. W. SANFORD  
 Mullard Space Science Laboratory, University College London, Holmbury St. Mary, Surrey

AND

R. D. JOSEPH AND W. P. S. MEIKLE  
 Blackett Laboratory, Imperial College of Science and Technology, London  
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### ABSTRACT

During 1978 April, May, and July, HDE 226868, the optical counterpart of Cygnus X-1, was repeatedly observed in the ultraviolet with the *IUE* satellite. Some X-ray and infrared observations have been made during the same period. The general shape of the spectrum is that expected from a late O supergiant. Strong absorption features are apparent in the ultraviolet, some of which have been identified. The equivalent widths of the most prominent lines appear to be modulated with the orbital phase. This modulation is discussed in terms of the ionization contours calculated by Hatchett and McCray, for a binary X-ray source in the stellar wind of the companion.

*Subject headings:* infrared: sources — line identifications — spectrophotometry — stars: individual — ultraviolet: spectra — X-rays: binaries

### 1. INTRODUCTION

Cygnus X-1, one of the strongest X-ray sources in the sky, was discovered by Bowyer *et al.* (1965). A radio counterpart was found by Braes and Miley (1971) and Hjellming and Wade (1971), following the report of fast variability by the *Uhuru* team (Oda *et al.* 1971). The identification in the optical is due to Webster and Murdin (1972) and Bolton (1972*a*), who showed that the 9 mag star HDE 226868, inside the radio error box, was a single-line binary with period 5<sup>d</sup>.6 and proposed that the invisible companion was a black hole.

High and low states with durations of years are observed in the X-rays with softer and harder spectra, respectively (e.g., Oda 1977). The transition is rather abrupt and probably associated with activity in the radio band. Absorption dips and troughs lasting  $\sim \frac{1}{2}$  hr are observed sporadically near orbital phase 0 (for a review see Oda 1977).

Walborn (1973) has classified HDE 226868 as an O9.7 Iab (p var). The blue-violet absorption spectrum is normal. A peculiar feature is the variable emission line He II ( $\lambda 4686$ ), the radial velocity curve of which is shifted by about 120° with respect to that of its absorption lines (see Bolton 1972*b*, 1975; Walker, Young, and Glaspey 1978, and references therein). This line is observed in several X-ray sources and in the case of Cyg X-1 is interpreted to arise from the gas streaming onto the invisible secondary (Smith *et al.* 1973). Absorption in this stream is probably responsible for the X-ray dips.

Light variations were studied in detail (see Walker and Rolland Quintanilla 1978, and references therein) showing a double maximum light curve and an amplitude in the B band of 0.04 mag. No direct evidence of the presence of an accretion disk is available, and in fact it is not clear whether accretion onto the secondary

TABLE 1  
BASIC PARAMETERS OF THE SYSTEM  
HDE 226868 = CYGNUS X-1

Parameter	Value
Spectral type	O9.7 Iab
$T_{\text{eff}}$	$3 \times 10^4$ K
$U$	9.83 mag
$B$	9.68 mag
$V$	8.87 mag
$J$ (1.24 $\mu\text{m}$ )	6.79 mag
$H$ (1.65 $\mu\text{m}$ )	6.63 mag
$K$ (2.22 $\mu\text{m}$ )	6.50 mag
Distance	1.8 Kpc
Radius	$18 R_0$
Period	5.6 days
$M$	$21 M_0$
$L$	$2.3 \times 10^5 L_0$
$M_x$	$13 M_0$
$L_x$	$1.8 \times 10^3 L_0$

Note.—The data are taken from Bregman *et al.* 1973 and Conti 1978. The infrared magnitudes are from our observations.

occurs through a disk or is spherical (see, e.g., Petterson 1978). Models of the X-ray emission have been constructed for both cases (for a review, see Lightman, Shapiro, and Rees 1978).

The basic parameters of the system are summarized in Table 1. For a detailed description of the Cyg X-1/HDE 226868 system, we refer to Oda (1977), Bahcall (1978), Conti (1978), and Eardley *et al.* (1978).

In this paper we report some UV observations of HDE 226868 taken with the *International Ultraviolet Explorer* (*IUE*; see § II), during the 2 weeks when the satellite was continuously dedicated to the study of

X-ray sources. X-ray and IR observations taken at the same time are described in §§ III and IV. The results are discussed in § V. Here the UV spectrum is compared with Kurucz's (1979) model atmospheres, and the effect of X-ray heating on the wind of the primary is considered.

## II. *IUE* OBSERVATIONS OF HDE 226868

### a) *The Observations*

A general description of *IUE* and the data acquisition is given by Boggess *et al.* (1978*a, b*). The 14 low-resolution ( $\sim 6$  Å) spectra we examine here were produced at the Goddard station. A Journal of Observations is given in Table 2. Spectrum SWP 1273 was taken during the commissioning phase of activity of the *IUE* (Dupree *et al.* 1978).

Spectra SWP 1979 and 2049 have been reprocessed in order to correct for errors of the intensity transfer functions, following the procedure reported in *IUE/ESA Newsletter*, No. 4. The orbital phases are those derived by Kemp (1977) from optical photometry, with  $P = 5.600000 \pm 0.00024$  days and epoch 1976 January 26.745 UT for inferior conjunction of the primary.

The mean of the calibration curves produced by US and UK project groups (Bohlin *et al.* 1980) was used. The reduced spectra are given in Figures 1 and 2. Spectra LWR 1428 and 1439 (Figs. 2*e* and 2*f*) were taken with the small aperture diaphragm. In order to compare these latter spectra with the others, the fluxes were multiplied by a factor of 2, which is the normal light loss for that observing mode. However, because of the uncertainty of the procedure, these spectra have not been used for photometric purposes.

TABLE 2  
JOURNAL OF OBSERVATIONS OF HDE 226868

Image	Aperture	Exposure (min)	Phase	Date of Observation (JD 2,440,000 +)	Integral Flux ( $10^{-10}$ ergs $\text{cm}^{-2}$ $\text{s}^{-1}$ )	
					1250–1900 Å (SWP)	1950–2500 Å (LWR)
SWP 1273...	large	40	0.93	3599.03	1.04	...
SWP 1445...	large	50	0.32	3629.26	1.16	...
LWR 1425...	large	15	0.43	3629.87	...	0.48
SWP 1451...	large	50	0.44	3629.92	1.13	...
SWP 1460...	large	45	0.62	3630.90	1.16	...
LWR 1428...	large	10	0.62	3630.93	...	0.46
LWR 1428...	small	10	0.63	3630.95	...	...
SWP 1478...	large	45	0.97	3632.89	...	...
LWR 1439...	large	10	0.98	3632.91	...	0.40
LWR 1439...	small	10	0.98	3632.93	...	...
SWP 1515...	large	45	0.05	3638.91	1.06	...
SWP 1979...	large	45	0.27	3701.76	1.14	...
LWR 1819...	large	9	0.28	3701.80	...	0.52
SWP 2049...	large	50	0.54	3708.86	1.12	...

Note.—SWP: short-wavelength range 1150–2000 Å; LWR: long-wavelength range 1850–3350 Å; small aperture: 3", circular; large aperture: 10"  $\times$  20", oval; orbital phases are taken from Kemp 1977.

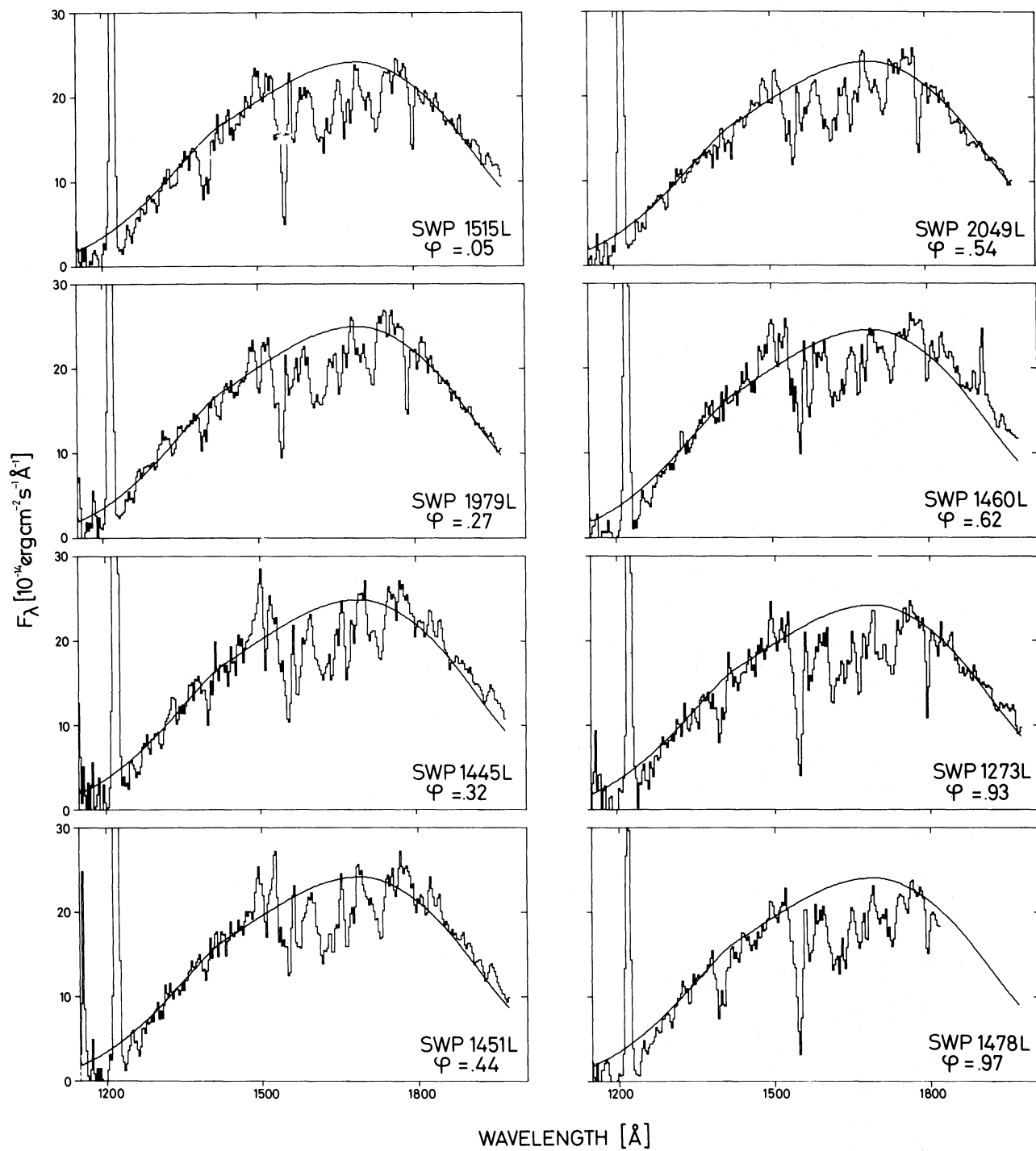


FIG. 1.—IUE spectra in the short-wavelength range. The solid lines represent reddened ( $E_{B-V}=1.06$ ) blackbody distributions with  $T=2.5 \times 10^4$  K. For each spectrum the image number and the orbital phase are given.

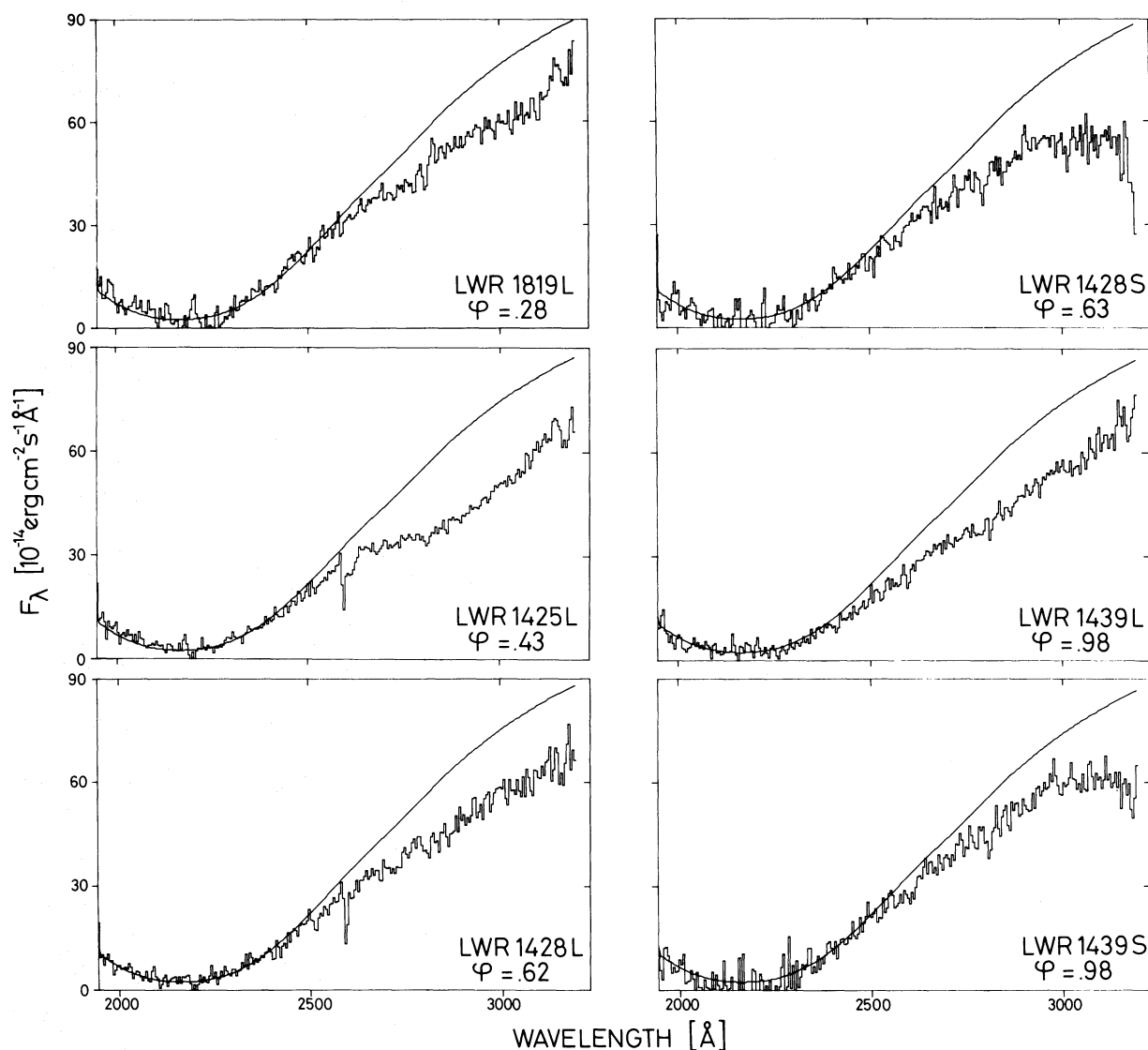


FIG. 2.—IUE spectra in the long-wavelength range. Spectra LWR 1428 and LWR 1439 have been taken with the small aperture and have been multiplied by 2 to compensate for light lost beyond the edges of the aperture. All spectra taken with the large aperture are partially saturated in the 2600–3200 Å region.

### b) The Continuum

An examination of the spectra indicates the difficulty of estimating the continuum since many strong absorption features are present. As a guideline, we plotted, for each phase, a blackbody distribution with the fixed temperature  $T=2.5 \times 10^4$  K (see Figs. 1 and 2), obtained normalizing to the B flux at the proper phase and fitting in two regions where absorptions are not prominent, namely in the intervals 1250–1350 Å, and 2350–2500 Å. The B light curve of Walker and Rolland Quintanilla (1978) (which exhibits a minimum

to maximum variation of 4%) was used with the phases given by Kemp (1977). The reddening curve of Seaton (1979) is assumed with  $E_{B-V}=1.06$ , which is intermediate between the values respectively given by Bregman *et al.* (1973) and Wu, Van Duinen, and Hammerschlag-Hensberge (1975). The adopted temperature obviously depends on the assumed  $E_{B-V}$ . For example  $E_{B-V}=1.02$  yields a temperature lower by 2000 K, while the temperature is increased by the same amount for  $E_{B-V}=1.12$ . This uncertainty is not really relevant, because the temperature derived in this way has no direct physical meaning, and the fitted

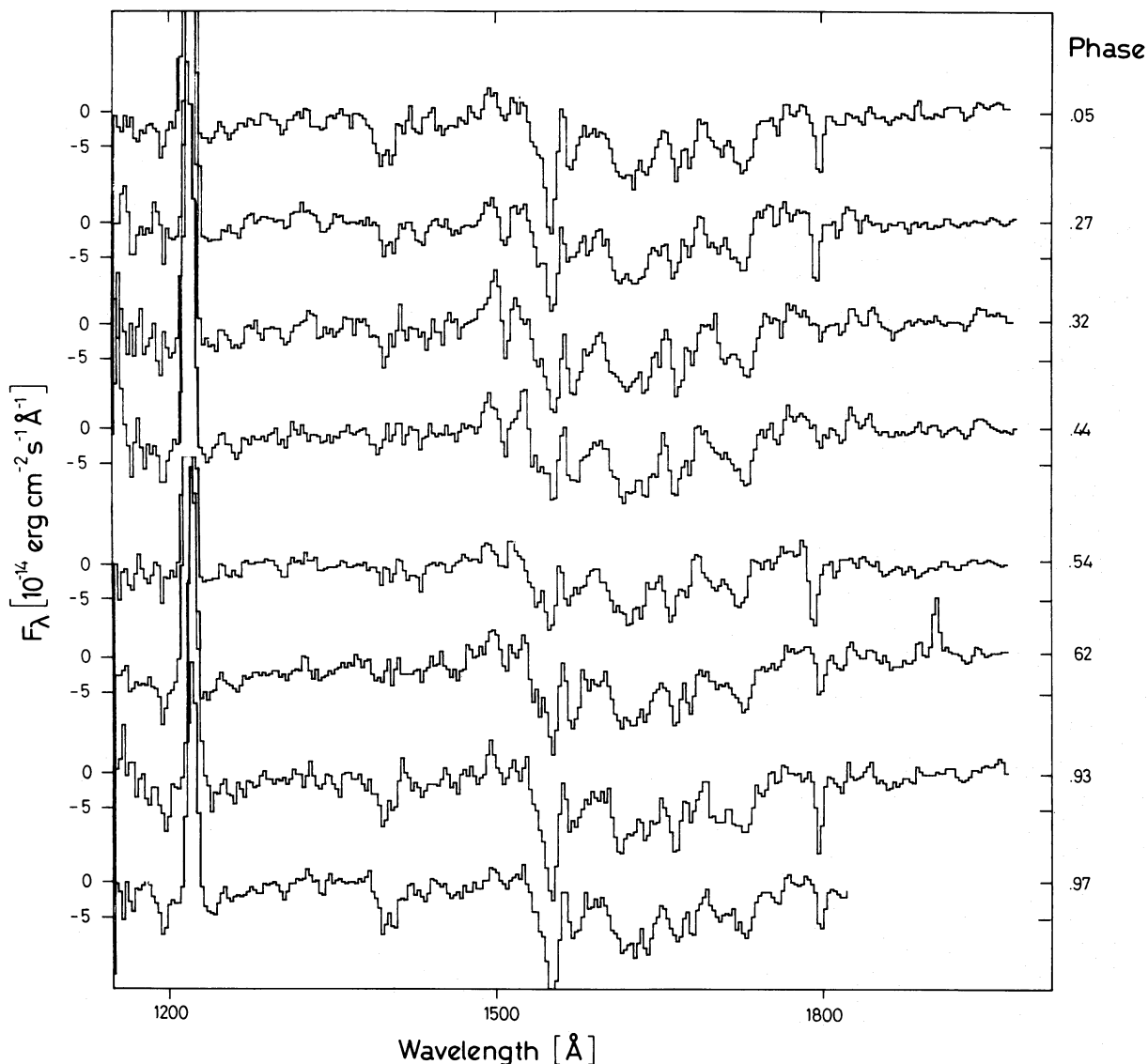


FIG. 3.—Rectified spectra in the short-wavelength range. For each spectrum the zero level together with the ordinate scale are given.

blackbody distributions are used only for determining unambiguously the equivalent widths of the absorption lines.

In Table 2 we list the integrated fluxes in the 1250–1900 Å and 1900–2500 Å bands for the SWP and LWR exposures. The rms variation of these fluxes is  $\pm 5\%$ —comparable to the  $\pm 3\%$  reported by Bohlin *et al.* (1979), as the fluctuation in *IUE* observations of standard stars. However, the average of the integrated fluxes  $F_{1250-1900}$  of the two short-wavelength spectra taken near phase 0 is  $1.05 \pm 0.01 \times 10^{-10}$  ergs  $\text{cm}^{-2}$   $\text{s}^{-1}$ , while the five spectra taken at other phases give a mean value of  $1.14 \pm 0.02$ . This suggests the existence of a minimum of the UV intensity around phase 0. The

maximum to minimum variability is  $8 \pm 2\%$ . A similar variability is found also in the (1950–2500 Å) range where, however, we have fewer observations.

### c) Lines

Figure 3 shows our short-wavelength rectified spectra, obtained by plotting the differences between the observed spectra and the blackbody curves of Figure 1. All spectra exhibit several absorption lines, some of which appear to be variable. In Table 3, the most prominent absorption features in the short-wavelength range are reported together with proposed identifications, recognizing that many of these are blends and will require high-resolution observations for positive

TABLE 3  
STRONG ABSORPTION FEATURES AND PROPOSED IDENTIFICATIONS

Proposed Identifications	Measured Wavelengths (in Å) on SWP 1515
Si II $\lambda$ 1261; $\lambda$ 1265	1258
O I $\lambda$ 1302; $\lambda$ 1306	1306
C II $\lambda$ 1334, $\lambda$ 1336	1340
Not identified	1369
Si IV $\lambda$ 1394, $\lambda$ 1403	1396, 1406
Not identified	1450
Not identified	1481
(Fe II $\lambda$ 1507 and Fe III $\lambda$ 1505)?	1509
C IV $\lambda$ 1548, $\lambda$ 1550	1550
Blend of Fe II lines	1566 to 1575
Blend	1614 to 1624
He II $\lambda$ 1640	1638
(O III $\lambda$ 1663 and Fe II $\lambda$ 1671)?	1666
Not identified	1679
Not identified	1707
N IV $\lambda$ 1719 and Al II $\lambda$ 1719 to $\lambda$ 1725	1725
Not identified	1760

identification. The measured wavelengths refer to spectrum SWP 1515. Shifts in wavelength in other spectra are within the expected accuracy of the wavelength scale, i.e., no greater than 5 Å. In the long-wavelength spectra, strong lines are absent, and because of the poor resolution, identification becomes hazardous. An absorption feature probably corresponding to Mg II ( $\lambda\lambda$ 2796,2803) appears in all six spectra.

In Table 4, equivalent widths of the most prominent features are reported for each spectrum, using two different estimates of the continuum. The former (left columns) refers to the blackbody continuum (see Fig. 1), the latter (right columns) to a continuum fitted by eye. In both cases integration intervals taken into account are determined by eye.

One can note that the equivalent widths of Si IV and C<sub>IV</sub> vary similarly with orbital phase, being largest around phase 0 (see Fig. 4). The intrinsic errors on each point are about 6% ( $2\sigma$  reproducibility—Bohlin *et al.* 1979), but the main source of uncertainty is the choice of the continuum. This can affect significantly the value of single determinations but, as is apparent from Figure 4, not the general behavior.

The phase modulation of the equivalent widths of some absorption lines can explain at least part of the variability of the integrated fluxes reported above.

### III. X-RAY OBSERVATIONS

Cyg X-1 was observed with the *Copernicus* satellite from JD 2,443,630.83 ( $\phi=0.62$ ) to JD 2,443,635.96 ( $\phi=0.53$ ), which covers *IUE* observations SWP 1460, SWP 1478, LWR 1428, and LWR 1439. A plot of the data is given in Figure 5. Each point represents approximately  $1\frac{1}{2}$  hr integration of the background-

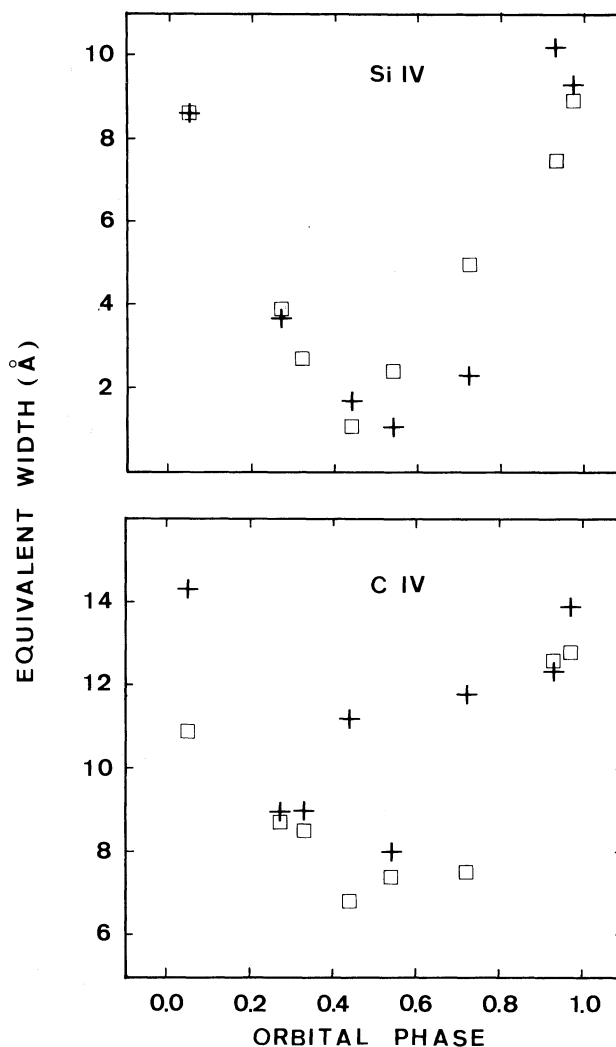


FIG. 4.—Equivalent widths of Si IV (upper) and C IV (lower) vs. orbital phase. Open squares correspond to the blackbody continuum, and crosses to a continuum fitted by eye.

corrected raw data (which has approximately 86 s sampling time). The error bars on the counts represents  $1\sigma$ . The detector energy range was 3.8–8.8 keV. The spectral hardness ratio (4.6–8.8 keV counts to 3.8–4.6 keV counts) is also shown as an indication of possible spectral changes. No significant change of the spectrum indicating increased photoelectric absorption was found, but just over 2 days into the observations, at  $\phi=0$ , there is a significant intensity drop which could be due to electron scattering. The observations show that the source was in a low state.

### IV. INFRARED PHOTOMETRY OF CYGNUS X-1

Cyg X-1 was observed between 04:30 UT and 05:30 UT on 1978 May 1, when the phase was 0.41. This is very close to *IUE* observation SWP 1415. The *J*, *H*,

TABLE 4  
EQUIVALENT WIDTHS OF SELECTED LINES

ION	SWP 1515 phase 0.05	SWP 1979 phase 0.27	SWP 1445 phase 0.32	SWP 1451 phase 0.44	SWP 2049 phase 0.54	SWP 1460 phase 0.62	SWP 1273 phase 0.93	SWP 1478 phase 0.97
C II $\lambda\lambda$ 1334,1336 and O IV 1338 .....	2.1	0.9	4.4	0.8	1.1	1.2	1.0	0.5
Si IV $\lambda\lambda$ 1394,1403 .....	8.6	3.9	...	1.1	2.4	0.5	7.5	8.9
C IV $\lambda\lambda$ 1548,1550 and Fe III $\lambda$ 1540 .....	10.9	8.7	9.0	6.8	7.4	7.5	12.6	12.8
Fe II $\lambda$ 1564 to $\lambda$ 1575 ...	3.3	3.1	3.5	3.1	3.3	3.7	4.1	4.3
O III $\lambda$ 1663 and Fe II $\lambda$ 1671 .....	6.6	4.5	3.5	5.5	5.6	5.5	7.6	7.5

Note.—Equivalent widths, in  $\text{\AA}$ , measured with reference to a blackbody continuum (left columns) and to a continuum fitted by eye (right columns).

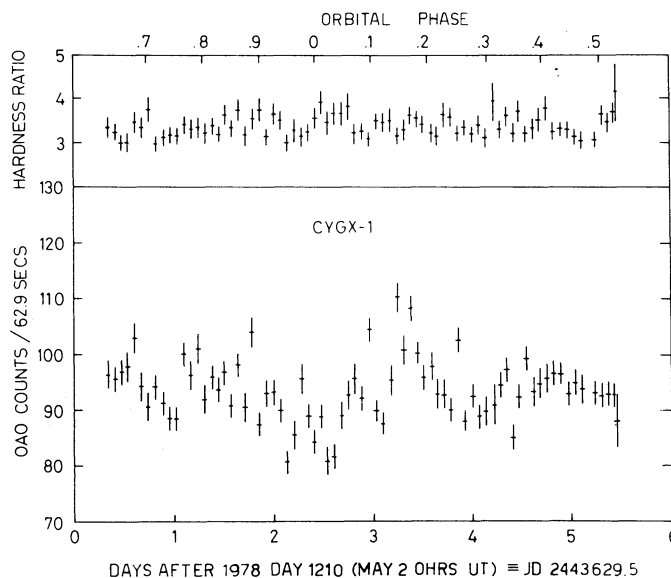


FIG. 5.—X-ray intensities of Cyg X-1; *Copernicus* counts and hardness ratios. The upper abscissa scale gives the orbital phase.

and  $K$  wavelengths were centered at  $1.24 \mu\text{m}$ ,  $1.65 \mu\text{m}$ , and  $2.22 \mu\text{m}$ , respectively. The calibration star was  $\alpha$  Bootis. The measured magnitudes are reported in Table 1. The  $1\sigma$  accuracy is 1%. The magnitudes were derived using the absolute calibration of Thomas, Hyland, and Robinson (1973) and Gehrz, Hackwell, and Jones (1974). Values of  $A_\lambda/E_{B-V}$  in the near-infrared were obtained from the studies of OBA stars by Schultz and Wiemer (1975), Whittet, von Breda, and Glass (1976), and Sneden *et al.* (1978). Our data, dereddened for  $E_{B-V}=1.06$  are consistent with a Rayleigh-Jeans spectrum and so may be interpreted as the long-wavelength tail of the photospheric emission spectrum of the O9.7 star.

For a distance modulus of 12 (Bregman *et al.* 1973), the absolute magnitude at  $2.22 \mu\text{m}$  becomes  $-5.9 \pm 0.11$ , i.e., about 50% brighter than that quoted by Becklin *et al.* (1972). However, they do not quote their adopted values for distance modulus or interstellar reddening, and so little significance can be attributed to this difference.

#### V. DISCUSSION

In Figure 6, the UV spectra SWP 1515 and LWR 1439 have been combined in order to compare the general shape of the UV spectrum around phase 0, when the primary is in front of the X-ray source, with the two model atmospheres which, in the grid calculated by Kurucz (1979), are closer to the spectral classification of the primary. Curve *a* represents Kurucz's model for  $T_{\text{eff}}=30,000$  K and surface gravity  $\log g=3.5$ , and curve *b* corresponds to  $T_{\text{eff}}=25,000$  and  $\log g=3.0$ .

In order to obtain a satisfactory agreement with data, a reddening  $E_{B-V}=1.12$  and  $E_{B-V}=1.06$ , respectively, for the two models has been adopted with the reddening curve by Nandy *et al.* (1976) for the Cygnus region. The comparison is limited to  $\lambda > 1400$  as the reddening curve is given only for these wavelengths.

The UV absorption features are typical of those observed in supergiants of similar spectral type (Panek and Savage 1976; Nandy 1977; Code 1977; Code and Meade 1979). In particular Panek and Savage (1976) give for supergiants between O9 and B0 equivalent widths of  $9 \text{ \AA}$  and of  $12 \text{ \AA}$  for Si IV ( $\lambda 1400$ ) and C IV ( $\lambda 1550$ ), respectively. Our measurements of Si IV at phase 0 are in good agreement with the above value, while at phase 0.5 the equivalent width is  $3 \text{ \AA}$ . C IV, which in our measurement may be contaminated by Fe III ( $\lambda 1540$ ), does not vary as strongly with phase. Its equivalent width around phase 0 is again close to that given by Panek and Savage (1976).

An interpretation of the phase dependence of the equivalent widths (Fig. 4) can be given in terms of the discussion by Hatchett and McCray (1977) of X-ray heating on stellar winds.

In particular they show that an X-ray source produces a region of high ionization which moves with the X-ray source. When this region is in front of the primary ( $\phi=0.5$ ), the intensity of the absorption lines is minimum, because of depletion of the absorbing ion species, while the opposite occurs at  $\phi=0$ .

In the Hatchett and McCray (1977) model the wind is taken as spherically symmetric around the primary, and the X-ray source is emitting isotropically. In the optically thin case the ionization structure is determined



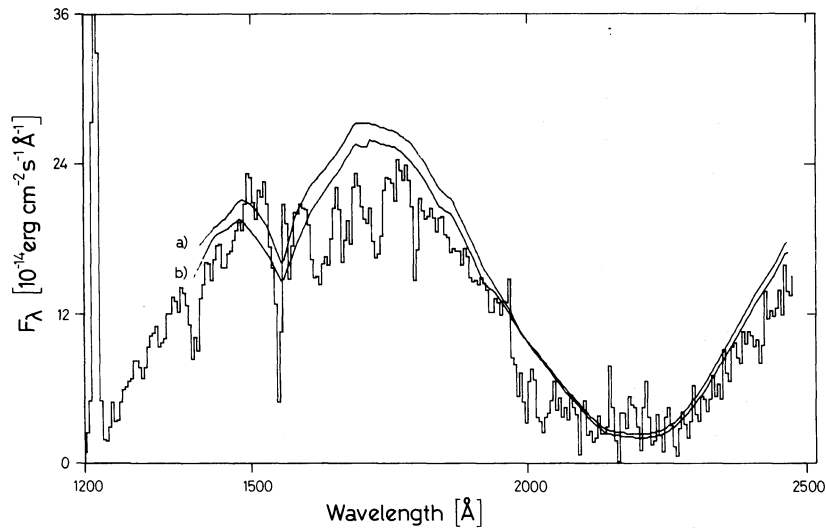


FIG. 6.—Comparison of the measured UV spectrum around phase 0 with Kurucz's model atmospheres. Curve *a* is for  $T_{\text{eff}}=30,000$  K,  $\log g=3.5$ , and  $E_{B-V}=1.12$ . Curve *b* is for  $T_{\text{eff}}=25,000$  K,  $\log g=3.0$ , and  $E_{B-V}=1.06$ .

by the single parameter  $\xi=L_X/n(r_x)^2$ , where  $L_X$  is the luminosity of the X-ray source,  $r_x$  is the distance from the X-ray source, and  $n$  is the number density. It may be useful to use the dimensionless parameter  $q=\xi/\xi_0$ , where  $\xi_0=L_X/n_0 d^2$ ,  $d$  is the distance of the center of the primary from the X-ray source, and  $n_x$  is the density at the X-ray source. Assuming a constant velocity wind, the loci of constant  $q$  are spheres with centers on the straight line connecting the centers of the two stars. Their radius is  $dq^{1/2}(1-q)^{-1}$ , and the centers are located on the side closer to the X-ray source for  $q<1$ , or closer to the primary of  $q>1$ . Hatchett and McCray (1977) show that the surfaces of the constant column density are in good approximation spheres congruent to those of constant  $q$ . Within this simplified picture one expects that the line intensity modulation is symmetric around phase 0.5 (see Fig. 4).

We now discuss the extension of the region transparent to the C IV and Si IV lines, following the paper by Dupree *et al.* (1980) on Vela X-1. First, one must estimate the total depth  $\tau$  of the wind. Our low-resolution UV spectra do not allow a direct estimate of  $\dot{M}$ . Taking  $\dot{M}\approx 2.5\times 10^{-6} M_\odot \text{ yr}^{-1}$ , as suggested by Hutchings (1976) on the basis of observations in the visible, and a wind velocity  $V\approx 1600 \text{ km s}^{-1}$  (Conti 1978) one derives  $\tau=4\times 10^{22} \text{ cm}^{-2}$ . The optical depths for C IV and Si IV absorptions are therefore  $\tau_{\text{C IV}}=8\times 10^3 X_{\text{C IV}}$  and  $\tau_{\text{Si IV}}=1.5\times 10^3 X_{\text{Si IV}}$ , where  $X_{\text{C IV}}$  and  $X_{\text{Si IV}}$  are the fraction of carbon and silicon in the given state of ionization. Therefore the regions where  $X_{\text{C IV}}<5\times 10^{-4}$  and  $X_{\text{Si IV}}<2.5\times 10^{-4}$  are certainly transparent to the wavelengths corresponding to the C IV and Si IV lines. From the calculations by Hatchett, Buff, and McCray (1976) and preliminary results by

Kallmann and McCray (1979) one knows that the above concentrations correspond to  $\log \xi=1.7$  and  $\log \xi=1.2$  for C IV and Si IV, respectively. Taking an X-ray luminosity  $L_X\approx 1.8\times 10^3 L_\odot$ , a separation  $d\approx 43 R_\odot$ , and  $n_0\approx 6\times 10^9 \text{ cm}^{-3}$ , one has  $q\approx 0.4$  for C IV and  $q\approx 0.1$  for Si IV. The values of  $q$  correspond to transparent regions lying both outside the conical shadow of the primary and also outside a radius of  $43 R_\odot$  for C IV and  $16 R_\odot$  for Si IV.

This result indicates that the region transparent to Si IV is more extended than that transparent to C IV, in agreement with our observations, which show that departure of the equivalent widths from that of an undisturbed star are larger for Si IV than for C IV. The shape of the opaque and transparent regions discussed above should be considered as a first approximation, since there are several reasons why the values of  $q$  can differ from the calculated ones. The largest uncertainty is in the adopted value of the wind parameters in the hypothesis that the wind is spherically symmetric around the primary without being influenced by the X-ray source. The gas streaming mentioned in the Introduction (§ I), connected with the X-ray absorption dips is in fact an indication of the roughness of the spherical approximation. Moreover, the Cyg X-1 X-ray spectrum is significantly different from the bremsstrahlung spectrum with  $kT=10 \text{ keV}$  considered by Hatchett, Buff, and McCray (1976).

If the overall picture is correct, following Hatchett and McCray (1977), we would expect the shape of P Cyg profiles to depend on the orbital phase. The investigation of this requires observations at high dispersion, as was done for the much brighter system Vela X-1=HD 77581 by Dupree *et al.* (1980).

Finally, we recall that the observations reported in this paper were taken during the low state of the X-ray source (see § III). Since the high state corresponds to a higher soft X-ray luminosity, one expects variability of temperature and density, of the plasma surrounding the primary, which should affect the UV spectral lines.

A comparison of the UV observations in the two states should therefore be of great interest.

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J. BLACK, R. J. DAVIS, A. K. DUPREE, H. GURSKY, L. W. HARTMANN, and J. RAYMOND: Harvard-Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138

L. CHIAPPETTI and A. TREVES: Istituto di Fisica dell'Università, Via Celoria 16, Milan, Italy

R. D. JOSEPH and W. P. S. MEIKLE: Blackett Laboratory, Imperial College of Science and Technology, Prince Consort Road, London SW7, England

T. A. MATILSKY: Department of Physics, Rutgers University, Piscataway, NJ 08854

G. POLLARD and P. W. SANFORD: Mullard Space Science Laboratory, University College London, Holmbury St. Mary, Surrey, England

E. G. TANZI: Laboratoria di Fisica Cosmica e Tecnologie Relative, Via Bassini 15, Milan, Italy

M. TARENGHI: European Southern Observatory, CERN, Geneva, Switzerland

P. VANDEN BOUT and F. SANNER: Physics Department, University of Texas, Austin, TX 78712